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ETH

Eidgenössische Technische Hochschule Zürich Swiss Federal Institute of Technology Zurich

Acoustics I: room acoustics

Kurt Heutschi 2022-12-12

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sound propagation in rooms



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approaches and indicators in room acoustics:

- atistical assumption: sound field is composed of direct sound and *diffuse field*. Indicators: reverberation time, direct/diffuse energy ratio.
- eometrical assumption: propagation along straight rays (high frequency approximation). Indicators: reverberation time, additional room acoustical parameters.
- wave based solutions of the wave equation, analytical solutions for simple geometries only, numerical methods for the general case.

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applicability of the approaches:

	small rooms	large rooms
low frequencies	W	(W),G,S
high frequencies	(W),G,S	G,S

- S: statistical
- ► G: geometrical
- ► W: wave based

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diffuse sound field

- statistical approach ignoring specific position of features
- ▶ fundamental assumption: direct + diffuse sound field
- properties of a diffuse sound field:
 - constant sound energy density
 - no dominant sound incidence direction
- conditions for a diffuse sound field in practice:
 - only little total absorption
 - homogeneous distribution of the absorption

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sound energy density and intensity on a surface

energy density and intensity

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- \blacktriangleright situation: room with a diffuse sound field of energy density w
- wanted: normal component of sound intensity on the surfaces
 solution:
 - formulation of the relation between a small volume element and a surface element
 - integration over the volume

energy density and intensity

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energy density and intensity

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energy *E* captured by dS, stemming from dV:

$$E = \mathrm{d}S\cos\theta\frac{w\cdot\mathrm{d}V}{4\pi r^2}$$

in spherical coordinates:
$$dV = r^2 \cdot dr \cdot d\theta \cdot \sin(\theta) \cdot d\phi$$

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energy density and intensity

- power W captured by dS corresponds to the energy collected in 1 second
- \blacktriangleright \rightarrow energy incident from a half-sphere with radius $R = c \cdot 1$ sec

$$W = \frac{w \mathrm{d}S}{4\pi} \int_{0}^{c \cdot 1 \sec} \int_{0}^{2\pi} \int_{0}^{\pi/2} \cos(\theta) \sin(\theta) \mathrm{d}\phi \mathrm{d}\theta \mathrm{d}r = \frac{wc}{4} \mathrm{d}S = I \mathrm{d}S$$

and for the intensity:

$$I_{\text{wall}} = rac{w \cdot c}{4}$$

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power balance in the diffuse field

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power balance in the diffuse field

- situation: sound source emits power W_{source} in a room, assuming a diffuse field
 - wanted: sound pressure in steady state
 - solution:
 - interpret diffuse field as superposition of plane waves
 - ▶ balance condition for $W_{\text{source}} = W_{\text{absor}}$

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power balance in the diffuse field

absorbed power:

where:

$$A = \sum_{i} \alpha_{i} \cdot S_{i}$$

 $W_{ ext{absor}} = I_{ ext{wall}} A = rac{WC}{A} A$

from $W_{\text{source}} = W_{\text{absor}}$ follows:

$$w = \frac{4W_{\text{source}}}{Ac}$$

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power balance in the diffuse field

for a plane wave in a cylindrical tube of cross sectional area 1 m² and height $= c \cdot 1$ sec:

 $I = \frac{p^2}{\rho c} = wc$

and finally:

$$p_{\rm diffuse}^2 = rac{4W_{
m source}
ho c}{A}$$

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distance dependency in the sound field

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distance dependency of p in the sound field

- distance dependency of p in the sound field for an omnidirectional point source:
 - superposition of p^2 of direct and diffuse sound

direct sound:

diffuse sound:

$$p_{
m direct}^2 = rac{W_{
m source}
ho c}{4\pi r^2}$$

$$p_{\rm diffuse}^2 = rac{4W_{
m source}
ho c}{A}$$

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distance dependency in the sound field summation:

$$p^2 = p_{ ext{direct}}^2 + p_{ ext{diffuse}}^2 = W_{ ext{source}}
ho c \left(rac{1}{4\pi r^2} + rac{4}{A}
ight)$$



distance dependency in the sound field

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• critical distance r_c where $p_{direct}^2 = p_{diffuse}^2$:

from
$$\frac{1}{4\pi r_c^2} = \frac{4}{A}$$
 follows $r_c = \sqrt{\frac{A}{16\pi}}$

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reverberation, reverberation time

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reverberation, reverberation time

observation of the non-stationary process after muting the source:
 power balance:

$$0 = W_{\mathsf{absor}}(t) + V rac{\mathrm{d} w(t)}{\mathrm{d} t}$$

with

*W*_{absor}: absorbed power by the room surfaces *V*: room volume *w*: energy density

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reverberation, reverberation time

with
$$W_{absor} = wcA/4 \rightarrow$$
 differential equation for $w(t)$:

$$0 = \frac{w(t)c}{4}A + V\frac{\mathrm{d}w(t)}{\mathrm{d}t}$$

hypothesis for the solution:

$$w(t) = w_0 e^{bt}$$

comparison of coefficients yields: $b = -\frac{cA}{4V}$

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reverberation, reverberation time

time dependency of w(t)
ightarrow reverberation :



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reverberation, reverberation time

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- reverberation time T:
 - time span for a drop of w to 10⁻⁶ of its initial value
 corresponds to -60 dB

from:

$$e^{-rac{cA}{4V}T}=10^{-6}$$

follows (Sabine formula):

$$T = \frac{-\ln(10^{-6})4V}{cA} = \frac{0.16V}{A}$$

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reverberation time according to Eyring

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Eyring's reverberation time

- model for sound propagation process:
 - sound energy particles propagate along straight lines
 - after a hit at a room surface:
 - specular reflection
 - $energy_{new} = energy_{old} \cdot (1 \alpha)$
 - α: average absorption coefficient
 - distance between two reflections: mean free path length ℓ
 - for rectangular rooms: $\ell = \frac{4V}{S}$
 - V: room volume
 - S: room surface

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Eyring's reverberation time

energy E of a particle after N reflections:

$$E(N) = E_0(1-\alpha)^N$$

E drops to 10^{-6} after M reflections with

$$M = \frac{\ln(10^{-6})}{\ln(1-\alpha)} = \frac{-13.8}{\ln(1-\alpha)}$$

M reflections correspond to a path length $L = M \cdot \ell$ corresponding to a time *T* (reverberation time):

$$T = \frac{M\ell}{c} = \frac{-13.8 \cdot 4V}{\ln(1-\alpha)cS} = \frac{0.16V}{-\ln(1-\alpha)S}$$

comparison of Eyring and Sabine

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- always: $T_{\text{Eyring}} < T_{\text{Sabine}}$
- ▶ for $\alpha \rightarrow 0$: $T_{\text{Eyring}} = T_{\text{Sabine}}$
- ▶ for $\alpha \rightarrow 1$: $T_{\text{Eyring}} \rightarrow 0$, $T_{\text{Sabine}} > 0$

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example: reverberation time calculation

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Sabine reverberation time calculation

- rectangular room 10m x 6m x 4m with
 - front and rear wall: plaster
 - one side wall: windows
 - one side wall: wooden boxes
 - floor: carpet
 - ceiling: plaster

absorption coefficients α [.]:

	125	250	500	1k	2k	4k
plaster	0.01	0.01	0.02	0.02	0.03	0.04
window	0.1	0.04	0.03	0.02	0.02	0.02
wooden boxes	0.5	0.25	0.15	0.05	0.05	0.1
carpet	0.05	0.08	0.2	0.3	0.35	0.4

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Sabine reverberation time calculation

absorption $A = \alpha_i \cdot S_i$ [m²] for each surface:

	125	250	500	1k	2k	4k
front/rear wall	0.5	0.5	1.0	1.0	1.4	1.9
side wall 1	4.0	1.6	1.2	0.8	0.8	0.8
side wall 2	20.0	10.0	6.0	2.0	2.0	4.0
floor	3.0	4.8	12.0	18.0	21.0	24.0
ceiling	0.6	0.6	1.2	1.2	1.8	2.4
TOTAL	28.1	17.5	21.4	23.0	27.0	33.1

Sabine reverberation time calculation: result

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	125	250	500	1k	2k	4k
RT [s]	1.37	2.20	1.80	1.67	1.42	1.16
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reasons for deviations between measured and calculated reverberation times:

additional air absorption

- effect can be included by more subtle formulas
- non-diffuse sound fields
 - inhomogeneous absorber distribution
 - coupled rooms

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- inhomogeneous absorber distribution:
 - high absorption concentrated at one surface
 - sound energy may oscillate between parallel surfaces with low absorption
 - \blacktriangleright \rightarrow effect: longer reverberation times than predicted by Sabine/Eyring
- examples of inhomogeneous absorption:
 - rooms with exclusive absorption by audience
 - gyms with absorbing ceilings

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- coupled rooms:
 - rooms with different damping, connected to each other by a small common opening
 - source and receiver in the room with higher damping
 - \blacktriangleright \rightarrow sagging reverberation curves
- examples of coupled rooms:
 - churches with attached chapels
 - auditorium and foyer (coupling by open doors)
 - coupled reverberation chambers (KKL)

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- model of sound propagation:
 - sound propagation in form of energy droplets that travel along straight lines (rays)
 - \blacktriangleright \rightarrow high frequency approximation (\approx light)
 - consequences:
 - \blacktriangleright consideration of energies only \rightarrow wave phenomena are ignored:
 - no diffraction
 - no interference
 - no resonances

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- consideration of absorption in air and at room surfaces
 reflection at surfaces:
 - specular reflection
 - at smooth surfaces
 - diffuse reflection (Lambert's law)
 - at depth structured (compared to λ) surfaces

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specular reflection:

▶ angle of incident ray = angle of reflected ray



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diffuse reflection:

• Lambert's law: $I(\phi) \sim \cos(\phi)$



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- result of an analysis with geometrical acoustics:
 - energy impulse response
 - energy as a function of time of arrival (impulse excitation)



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objective room acoustical criteria

- energy impulse response allows for evaluation of a wide variety of room acoustical criteria such as:
 - reverberation time T
 - Early Decay Time EDT
 - clarity C80
 - Deutlichkeit D50
 - center time TS
 - strength G
 - lateral energy fraction LF

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objective room acoustical criteria

▶ reverberation time *T* [s]

- global descriptor
- ▶ just audible difference: 5%
- evaluation: inverse integration of the energy impulse response,
 - $T = 2 \cdot T'$ with T': time for drop from -5...-35 dB

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objective room acoustical criteria

► Early Decay Time *EDT* [s]

- subjectively more relevant than T
- ▶ just audible difference: 5%
- evaluation: inverse integration of the energy impulse response, $EDT = 6 \cdot T'$ with T': time for drop from 0...-10 dB

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clarity C80 [dB]

- ▶ just audible difference: 0.5 dB
- typical value for music rooms: 0 dB
- evaluation:

$$\mathcal{C}80 = 10 \log egin{pmatrix} 80 ms \ \int & h^2(t) \mathrm{d}t \ rac{0}{\infty} & \int & h^2(t) \mathrm{d}t \ \int & h^2(t) \mathrm{d}t \ \end{pmatrix}$$

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objective room acoustical criteria

Deutlichkeit D50 [%]

- ▶ just audible difference: 5%
- e.g. $D50 = 40\% \rightarrow 87\%$ intelligibility of syllables
- evaluation:

$$D50 = rac{\int\limits_{0}^{50ms} h^2(t) \mathrm{d}t}{\int\limits_{0}^{\infty} h^2(t) \mathrm{d}t} imes 100\%$$

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center time TS [ms]

▶ just audible difference: 10 ms

-

evaluation:

$$TS = rac{\int\limits_{0}^{\infty} t \cdot h^2(t) \mathrm{d}t}{\int\limits_{0}^{\infty} h^2(t) \mathrm{d}t}$$

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strength G [dB]

- just audible difference: 1 dB
- evaluation:

where

 $h_{f,10m}$: energy impulse response in 10 m under free field conditions.

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objective room acoustical criteria

- ▶ lateral energy fraction *LF* [%]
 - just audible difference: 5%
 - evaluation:



where

 h_{∞} : energy impulse response with a figure-of-eight microphone (lateral sensitivity).

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objective room acoustical criteria

optimal reverberation times:



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optimal values for large music rooms:

criterium	optimal range
EDT	1.82.2 s
C80	-2+2 dB
G	$> 0 \ dB$
LF	10% 35%

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listening examples

- general: T = 2 s
- varying values of other criteria

guitar ABCD

evaluation of transparency of examples A to D

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objective room acoustical criteria

values of the criteria: guitar 2128

Nr.	C80	ΤS	EDT	Т
1	-4.6	200	2.43	2.0
2	0.7	121	2.25	2.0
8	5.3	69	1.75	2.0

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room acoustical design criteria for large rooms

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design criteria for large rooms

- quietness
 - no audible noise
 - sufficient sound insulation to outside sources
 - sufficient sound insulation to adjacent rooms
 - sufficient low level of technical noise (ventilation, ...)
- direct sound
 - provide sufficient early energy to audience area
 - early energy: up to 50..80 ms after direct sound arrival
- reverberation
 - appropriate reverberation time in dependency of usage, room volume and room type

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design criteria for large rooms

lateral reflections

- feeling of spaciousness is created by signal differences at the two ears
 - Iow interaural cross correlation
 - need for strong lateral reflections
- diffusivity
 - temporally smeared reflections are beneficial
 - avoid discrete echoes, no focusing effects

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design criteria for large rooms

balance

- homogeneous signal strength distribution of all instruments in the audience area
- audibility on stage
 - musicians should hear each other reasonably well

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- statistical room acoustics:
 - calculation of reverberation times
- geometrical room acoustics:
 - construction of sound rays by hand
 - computer based ray tracing
 - computer based mirror source construction
- wave theoretical room acoustics:
 - calculation of room modes
 - numerical simulation (FDTD, BE, FE)
- experimental investigations:
 - in real rooms
 - is physical scale models
- auralisation:
 - listen to a source in a virtual room

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tools for the acoustical design of rooms

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statistical room acoustics: reverberation time calculations

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Statistical room acoustics: reverberation time calculations

- formulas of Sabine or Eyring
 - required data:
 - ▶ room volume V
 - area S and absorption coefficient α of the materials present in the room
- approximations for situations with dominating audience absorption (area S_a):

$$T \approx \frac{0.15 V}{S_a}$$

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geometrical room acoustics: exemplary sound rays
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geometrical room acoustics: exemplary sound rays

evaluation of basic forms

- shape of ceiling
- primary form of ground plan
- focusing elements
- orientation of reflector elements

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computer-assisted geometrical room acoustics: ray tracing

geometrical room acoustics: ray tracing

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- strategy:
 - source point emits energy particles
 - propagation along straight lines under consideration of specular or diffuse reflection and absorption
 - collection of rays that pass a sphere around each receiver
 - collection \rightarrow energy impulse response
- advantages:
 - diffuse reflections can be handled
 computational effort ~ Naturation
- disadvantages:
 - radius of receiver spheres > 0 \rightarrow fuzzyness in time and space
 no explicit truncation criterion

geometrical room acoustics: ray tracing

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geometrical room acoustics: ray tracing

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computer-assisted geometrical room acoustics: mirror source method

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geometrical room acoustics: mirror source method

- strategy:
 - construction of all mirror sources up to order n for all receivers
 - determine the "visible" (valid) mirror sources
 - apply weighting of $1/r^2$ and $(1-\alpha)$
 - ► collection: → energy impulse response
- advantages:
 - any desired temporal resolution possible
 clearly defined truncation criterion
- disadvantages:
 - diffuse reflections can't be handled
 - \sim computational effort $\sim \mathit{N}_{\mathsf{sources}} \cdot \mathit{N}_{\mathsf{receivers}}$

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scale model experiments

- investigation of sound propagation in physical scale models with help of measurements
- propagation effects scale with wavelength:
 - scaling of geometrical dimensions by 1/s
 - \blacktriangleright \rightarrow scaling of frequencies by s
- challenges:
 - identification of model materials with desired absorption properties
 - handling of extremely strong air absorption in the transformed frequency range
- ▶ typical scaling factors: 1:10 ... 1:50

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scale model experiments

example: Yokohama Minatomirai Concert Hall



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scale model experiments

example: Yokohama Minatomirai Concert Hall



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auralisation

- auralisation = provide listening impression of the acoustics of a virtual room
- listening impression is "ultimate" criterion
- very helpful in discussions with non-experts
- > audio material: dry recordings with no reverberation
 - example

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auralisation with scale models

- auralisation based on physical scale models:
 - frequency up-shift of the signal of interest
 - playback and recording in the model
 - frequency down-shift to create the signal to listen to

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auralisation with statistical room acoustics

- auralisation based on statistical room acoustics approach:
 - tunable parameters:
 - reverberation time
 - direct sound / diffuse sound ratio
 - reproduction: summation of original signal and output of reverberation engine

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auralisation based on geometrical room acoustics approach:

auralisation with geometrical room acoustics

- determination of individual energy impulse responses for different incidence directions
- reproduction
 - by cloud of loudspeakers
 - by headphones after filtering with head related transfer functions

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auralisation with geometrical room acoustics

example: time domain representation of HRTF sound incidence front (elevation = 0° , azimuth = 0°)



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auralisation with geometrical room acoustics

example: time domain representation of HRTF sound incidence right (elevation = 0° , azimuth = 90°)



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auralisation with geometrical room acoustics

example: time domain representation of HRTF sound incidence front up (elev. = 45° , azi. = 0°)



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additional effects usually not considered with statistical and geometrical room acoustics

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grazing incidence over audience areas

grazing incidence over audience areas

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- ► sound propagation at grazing incidence → strong excess attenuation due to:
 - destructive interference between direct sound and sound scattered at the bodies (heads and shoulders)

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grazing incidence over audience areas

example: measured levels at a receiver in the 12th row for a loudspeaker at a height of 1.2 $\scriptstyle ..$ 2.0 m



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reflection at finite surfaces

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reflection at finite surfaces

- geometrical interpretation of a reflection ignores:
 - wave covers a certain area for full reflection (frequency dependent)
- quantitative investigation by applying concept of Fresnel zones
- \blacktriangleright small surfaces \rightarrow scattering

reflection at finite surfaces

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• estimation of the lower limiting frequency f_u for full reflection:

$$f_u = \frac{2c}{(I\cos\beta)^2} \frac{d_{QR}d_{RE}}{(d_{QR} + d_{RE})}$$

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design of diffusely reflecting surfaces

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how to design a diffuser?

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- elements with diffuse/scattering reflection characteristics
- necessary condition:
 - introduction of a locally random phase response
 - incoherent superposition of reflection contribution of each surface element of the reflector
 - energetic integration for all directions



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- random phase response can be realized by
 - varying material properties
 - depth structured surface
 - channels of varying depth



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- variation of the channel depth according to:
 - white noise
 - sequence s of squared residues (Schroeder Diffuser):
 - $s = m^2 \mod N$ mit m = 0, 1, 2, ...; N: prime
 - ► s will become periodic in N
- frequency range for diffuse reflection:
 - $\lambda_{\max} \approx 2 \cdot \max.$ channel.depth • $\lambda_{\min} \approx 2 \cdot \max.$ channel.width
- \blacktriangleright caution: significant absorption due to large pressure differences from channel to channel \rightarrow pressure equalization without radiation

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Schroeder diffusor fractal extension:


diffusers: practical constructions

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reflection at convex and concave structures

reflection at convex and concave structures

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 $convex \rightarrow scattering$

 $concave \rightarrow focusing$

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convex structures are often beneficial

reflection at convex and concave structures



example: KKL

concave structures are often problematic!



pressure in circular room

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reflection at conic section geometries

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concave reflection at conic section geometries (source in focal point)



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reflection at segments of a circle

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reflection at concave segments of a circle

dome structures can often be approximated by

- ▶ parts of a sphere (3D)
- ▶ parts of circles (2D)

Cinderella (Kreis-Reflexionen.cdy)

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reflection at concave segments of a circle

segment of a circle appears (depending on source position) as:

- part of a hyperbola
- ▶ parabola
- ellipse

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segment of a circle appears (depending on source position) as:part of a hyperbola

- parabola
- ellipse

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reflection at concave segment of a circle

hyperbolic	parabolic	elliptic	elliptic
$x_Q < 0.5 R$ $x_{F2} = rac{x_Q}{2x_Q - 1}$	$x_Q = 0.5R$ $x_{F2} = \infty$	$\begin{array}{l} x_Q > 0.5R \\ x_{F2} = \frac{x_Q}{2x_Q - 1} \end{array}$	$y_Q^2 = x_Q - x_Q^2$ $y_{F2} = -y_Q$
divergent	parallel	focusing	focusing

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reflection at concave segment of a circle



reflection at concave segment of a circle

parabolic



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X_Q	<	0.5 <i>R</i>
X _{F2}	=	$\frac{x_Q}{2x_Q-1}$

divergent

 $\begin{array}{ll} x_Q = 0.5R & x_Q > 0 \\ x_{F2} = \infty & x_{F2} = \frac{1}{2} \end{array}$

parallel

focusing

focusing

reflection at concave segment of a circle



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reflection at concave segment of a circle

elliptic



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reflection at concave segments of a circle

study of two arrangements with interesting parabolic behavior:



receiver in focal point



source in focal point

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receiver in focal point: sound mirrors

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- large mirrors in form of parts of a sphere for sound amplification and localisation
- erected at the coast of England at the end of world war I
- goal: early detection of foreign airplanes by their acoustical signature
 acoustical aspects:
 - ▶ air absorption \rightarrow only low frequencies detectable in large distances ▶ focusing needed for $\lambda > 1$ m
 - ightarrow diameter of a mirror structure pprox several meters
- resulting performance: detection and localisation up to 20 km!

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sound mirrors

large mirrors in form of parts of a sphere for sound amplification and localisation

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sound mirrors

adjustment of focusing direction by moving the listening point in a spherically shaped mirror:





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other localisation methods

binaural localisation:



1921: military base near Washington, D.C.
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source in focal point: chapel in Guarda-Giarsun

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chapel in Guarda-Giarsun



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Guarda-Giarsun: interior views





Guarda-Giarsun: ground plan 1(2)

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source located in focal point, receiver distance 4 m

File: C:\KHE\MLSSA\GUAR-08.TIM 8-4-99 10:57 AM Impulse Response - volts

Guarda-Giarsun: impulse response measurement



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Guarda-Giarsun: amplification

source located in focal point, amplification as function of distance



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sound field in a circle

energy distribution in a circle 1(2)

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$\rightarrow \mathsf{Kreis}\text{-}\mathsf{Strahlen}.\mathsf{exe}$

energy distribution in a circle 2(2)

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 $ray(\phi) \rightarrow$ energy concentration on a circle with radius $r_k(\phi) = a \sin(\phi)$

$$ightarrow$$
 density: $rac{1}{rac{\partial r_k}{\partial \phi}} = rac{1}{a \cos(\phi)}$ max. for $\phi = 90
ightarrow r_k = a$

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wave fronts in circular rooms

development of wave fronts in a circular room:

movie: wave fronts in a circle

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wave fronts in circular rooms



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wave fronts in rectangular rooms

development of wave fronts in a rectangular room:

movie: wave front in a rectangular room

wave fronts in rectangular rooms

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chaotic behavior in stadium rooms

development of wave fronts in a stadium room:

movie: wave front in a stadium room

chaotic behavior in stadium rooms

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examples of buildings with round structures

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KIZA, church in Thun-Allmendingen

KIZA Thun-Allmendingen: elliptic ground plan

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KIZA Thun-Allmendingen: focusing effects in ellipses



Athanasius Kircher, 1684

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KIZA Thun-Allmendingen: intervention by a shield



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St. Pauls Cathedral, London

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St. Pauls Cathedral, London (1:3)

accessible gallery below the dome allows for communication by whispering over large distances.



St. Pauls Cathedral, London (2:3)

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St. Pauls Cathedral, London (3:3)

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Pantheon Rome 27 BC, Agrippa 118-128, Hadrian

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Pantheon Rome



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Pantheon Rome





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Plenarsaal Bonn

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Plenarsaal Bonn 4(6)

- remedial measure 1: all non transparent surfaces were covered with absorbing material
 - highly absorbing eagle wall
 - carpets along the walking zones
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Plenarsaal Bonn 5(6)

remedial measure 2: installation of transparent micro-perforated glass absorbers





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Plenarsaal Bonn 6(6)

remedial measure 3: guidance of sound reflections by tilted glass areas



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Forum im Pfalzkeller St. Gallen 1999, Santiago Calatrava

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Forum im Pfalzkeller St. Gallen: entrance



Forum im Pfalzkeller St. Gallen: groundplan

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Forum im Pfalzkeller St. Gallen: focusing problem

impulse response measurements:



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Forum im Pfalzkeller St. Gallen: focusing problems - remedial measures:

- central part of the ceiling was covered with 40 mm absorption
- installation of absorption along the curved rear wall
- installation of furniture with sound absorbing covers

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wave theoretical room acoustics

- in rooms with dimensions \approx wavelength of interest:
 - no longer free propagation
 - formation of resonances
 - wave based approach necessary
- strategy: search for the pressure field that fulfills:
 - wave equation
 - boundary conditions

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wave equation (Helmholtz equation):

$$\triangle \check{p} + \frac{\omega^2}{c^2} \check{p} = 0$$

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formulation of the boundary conditions (locally reacting surfaces):

$$\frac{p}{v_n} = Z$$

with Newton:

$$\frac{\partial p}{\partial n} = -\rho_0 \frac{\partial v_n}{\partial t}$$

follows an impedance relation in pressure only:

$$\frac{1}{\rho_0}\frac{\partial p}{\partial n} = -\frac{1}{Z}\frac{\partial p}{\partial t}$$

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boundary condition in complex writing for sinusoidal time dependencies:

 $\frac{1}{\rho_0}\frac{\partial \check{p}}{\partial n} = -\frac{1}{Z}\check{p}j\omega$

resp.

$$Z\frac{\partial \check{p}}{\partial n} + j\omega\rho_0\check{p} = 0$$

• special case: $Z \to \infty \Rightarrow \frac{\partial \check{p}}{\partial n} = 0$

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analytical solution for a rectangular room with rigid boundaries

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$$\begin{aligned} \frac{\partial \check{p}}{\partial x} &= 0 & \text{for } x = 0, x = L_x \\ \frac{\partial \check{p}}{\partial y} &= 0 & \text{for } y = 0, y = L_y \\ \frac{\partial \check{p}}{\partial z} &= 0 & \text{for } z = 0, z = L_z \end{aligned}$$

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analytical solution rectangular room

- sound field p̃(x, y, z) has to fulfill the boundary conditions and the Helmholtz equation
- ► Helmholtz equation in Cartesian coordinates:

$$\frac{\partial^2 \check{p}}{\partial x^2} + \frac{\partial^2 \check{p}}{\partial y^2} + \frac{\partial^2 \check{p}}{\partial z^2} + k^2 \check{p} = 0$$

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assumption for sound pressure $\check{p}(x, y, z)$:

$$\check{p}(x, y, z) = C \cos\left(\frac{n_x \pi x}{L_x}\right) \cos\left(\frac{n_y \pi y}{L_y}\right) \cos\left(\frac{n_z \pi z}{L_z}\right)$$

with

- n_x, n_y, n_z : positive integers including 0
- C: arbitrary constant
 - standing wave with nodes and anti-nodes (maxima)

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$$\check{p}(x, y, z) = C \cos\left(\frac{n_x \pi x}{L_x}\right) \cos\left(\frac{n_y \pi y}{L_y}\right) \cos\left(\frac{n_z \pi z}{L_z}\right)$$

check of fulfillment of boundary conditions (shown for *x*-coordinate):

$$\frac{\partial \check{p}}{\partial x} = -C \frac{n_x \pi}{L_x} \sin\left(\frac{n_x \pi x}{L_x}\right) \cos\left(\frac{n_y \pi y}{L_y}\right) \cos\left(\frac{n_z \pi z}{L_z}\right)$$
$$\rightarrow \frac{\partial \check{p}}{\partial x} = 0 \text{ at } x = 0, x = L_x \text{ for } n_x \text{ integer}$$

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- check of fulfillment of Helmholtz equation:
- ightarrow list of the second order derivatives with respect to space coordinates:

$$\frac{\partial^{2} \check{p}}{\partial x^{2}} = -C \frac{n_{x}^{2} \pi^{2}}{L_{x}^{2}} \cos\left(\frac{n_{x} \pi x}{L_{x}}\right) \cos\left(\frac{n_{y} \pi y}{L_{y}}\right) \cos\left(\frac{n_{z} \pi z}{L_{z}}\right)$$
$$\frac{\partial^{2} \check{p}}{\partial y^{2}} = -C \frac{n_{y}^{2} \pi^{2}}{L_{y}^{2}} \cos\left(\frac{n_{x} \pi x}{L_{x}}\right) \cos\left(\frac{n_{y} \pi y}{L_{y}}\right) \cos\left(\frac{n_{z} \pi z}{L_{z}}\right)$$
$$\frac{\partial^{2} \check{p}}{\partial z^{2}} = -C \frac{n_{z}^{2} \pi^{2}}{L_{z}^{2}} \cos\left(\frac{n_{x} \pi x}{L_{x}}\right) \cos\left(\frac{n_{y} \pi y}{L_{y}}\right) \cos\left(\frac{n_{z} \pi z}{L_{z}}\right)$$

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 \rightarrow inserted:

 $\begin{bmatrix} -\frac{n_x^2 \pi^2}{L_x^2} - \frac{n_y^2 \pi^2}{L_y^2} - \frac{n_z^2 \pi^2}{L_z^2} \end{bmatrix} \cdot \cos\left(\frac{n_x \pi x}{L_x}\right) \cos\left(\frac{n_y \pi y}{L_y}\right) \cos\left(\frac{n_z \pi z}{L_z}\right) + k^2 \cos\left(\frac{n_x \pi x}{L_x}\right) \cos\left(\frac{n_y \pi y}{L_y}\right) \cos\left(\frac{n_z \pi z}{L_z}\right) = 0$

o.k. for

$$k^2 = rac{n_x^2 \pi^2}{L_x^2} + rac{n_y^2 \pi^2}{L_y^2} + rac{n_z^2 \pi^2}{L_z^2}$$

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- Helmholtz equation in rectangular rooms with rigid boundaries is fulfilled for discrete wave numbers k only!
 - each positive (including 0) integer triple n_x, n_y, n_z defines an eigenvalue k_{nx}, n_y, n_z
 - the corresponding sound field $\check{p}(x, y, z)$ is a mode
- eigenvalue k can be expressed as frequency f:

$$f = \frac{c}{2} \sqrt{\frac{n_x^2}{L_x^2} + \frac{n_y^2}{L_y^2} + \frac{n_z^2}{L_z^2}}$$

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exemplary mode < 100 >:



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exemplary mode < 110 >:



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exemplary mode <200>:



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exemplary mode $<\!210\!>$:



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exemplary mode <320>:



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mode $<\!850\!>$?

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exemplary mode <850>:



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analytical solution rectangular room

- general findings for the pressure field:
 - all modes: pressure maximum in corners (0-dim)
 - modes with one n = 0: pressure maximum at corresponding edges (1-dim)
 - modes with two n = 0: pressure maximum at corresponding surfaces (2-dim)
- relevant for the positioning of pressure sensitive absorbers

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lowest eigenfrequencies for a small studio room of $8.5 \times 6.2 \times 2.6$ m (sorted in ascending order):

eigenfrequency [Hz]	n_{x}	n_y	n _z
20.0	1	0	0
27.4	0	1	0
33.9	1	1	0
40.0	2	0	0
48.5	2	1	0
54.8	0	2	0
58.4	1	2	0
65.4	0	0	1

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eigenfrequency differences between adjacent modes:

small studio room (8.5m, 6.2m, 2.6m, 137m³)



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eigenfrequency differences between adjacent modes:

Tonhalle (29.6m, 19.5m, 14m, 8080m³)



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estimate of the number of eigenfrequencies N_f from 0 up to f Hz in a room of volume V:

$$N_f pprox rac{4\pi}{3} V \left(rac{f}{c}
ight)^3$$

mode density dN_f/df (number of eigenfrequencies per Hz bandwidth):

$$\frac{dN_f}{df}(f) \approx 4\pi V\left(\frac{f^2}{c^3}\right)$$

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- damping in real rooms → finite quality factor of the resonances
 typical bandwidth: 1 Hz
- consideration of modes is important if they occur as isolated peaks (evaluated at -3dB)
- statistical and geometrical room acoustics applicable if mode peaks overlap
- "limiting frequency" = Schröder frequency for a mode density = 1/Hz
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lower limiting frequency for 1 mode per Hz bandwidth:



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source-receiver transfer function:

- considerations so far: possible sound fields in a room
- open question: can specific modes be excited by a volume source?
- ightarrow ightarrow source-receiver transfer function
 - solution of the wave equation with a source term

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sound pressure amplitude $\check{p}(E, \omega)$ at receiver E for excitation by volume source at Q:

$$\check{p}(E,\omega) \sim \omega \sum_{n} \frac{\check{p}_{n}(E)\check{p}_{n}(Q)}{(\omega^{2}-\omega_{n}^{2})K_{n}}$$

with

\sum_{n} : sum over all modes

 $\check{p}_n(E)$: complex sound pressure amplitude for mode *n* evaluated at point *E* $\check{p}_n(Q)$: complex sound pressure amplitude for mode *n* evaluated at point *Q* ω_n : mode specific eigenfrequency

 K_n : mode specific constant

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$$\check{p}(E,\omega) \sim \omega \sum_{n} \frac{\check{p}_{n}(E)\check{p}_{n}(Q)}{(\omega^{2}-\omega_{n}^{2})K_{n}}$$

- excitation of a mode is possible only if the volume source is located close to a pressure maximum
- excitation of all modes: source position in a corner
- plausibility check:

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$$\check{p}(E,\omega) \sim \omega \sum_{n} \frac{\check{p}_{n}(E)\check{p}_{n}(Q)}{(\omega^{2}-\omega_{n}^{2})K_{n}}$$

- excitation of a mode is possible only if the volume source is located close to a mode pressure maximum
- excitation of all modes: position in a corner
- plausibility check:
 - pressure maximum: part. velocity is small $\Rightarrow Z = \frac{p}{v} = Z_{rad} \rightarrow \infty$
 - ▶ pressure minimum: part. velocity is large $\Rightarrow Z = \frac{p}{v} = Z_{rad} \rightarrow 0$

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room acoustical design of small rooms

possible problems at *low frequencies*

- modal density is low, reverberation time is high
 - large pressure variations (in frequency and space)
 - bad transient behavior
- remedy:
 - damping of low frequencies by installation of resonance absorbers (membrane absorber, Helmholtz absorber, active absorber)

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example: listening room in a recording studio before installation of absorbers \rightarrow impulse response from loudspeaker to listener position:



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example: listening room in a recording studio before installation of absorbers \rightarrow amplitude response at listener position:



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resonance peak width Δf of a mode \leftrightarrow reverberation time *RT*:



 Δf : bandwidth at -3 dB points

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well conditioned studio room:

 $\Delta f \ge 5 \; {
m Hz} \qquad
ightarrow \; {
m RT} < 0.4 \; {
m s}$

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possible problems at mid and high frequencies

- too reverberant
 - blurred sound pattern
 - lack of transparency
 - flutter echoes
- early reflections
 - comb-filter effects, frequency response errors
 - stereo image errors

remedy:

- mid- and high frequency absorption by installation of porous absorbers
- sound propagation guidance by suitable reflector surfaces
- diffusely reflecting surfaces

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- listening room concept LEDE
 - Live-End-Dead-End
 - suppression of early reflections by absorption (dead end)



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- listening room concept RFZ
 - ► *R*eflection-*F*ree-*Z*one
 - suppression of early reflections by geometrical means



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reverberation time measurements

- global parameter to characterize a room (traditional parameter)
 averaging over several source and receiver positions
 measurement procedure:
 - interrupted noise or
 - inverse integration of squared impulse response
- important: documentation of air humidity and temperature (relevant at high frequencies)
- evaluation in octave bands
- decay curve is approximated by straight line (minimal least square deviation)

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- reverberation time measurement:
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- reverberation time measurement:
 - global parameter to characterize a room (traditional parameter)
 - averaging over several source and receiver positions
 - measurement procedure:
 - interrupted noise or
 - inverse integration of squared impulse response
 - important: documentation of air humidity and temperature (relevant at high frequencies)
 - evaluation in octave bands
 - decay curve is approximated by straight line (minimal least square deviation)

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impulse response measurement

- complete description of the propagation between source and receiver
 contains more information → evaluation of additional quantities such as EDT, C80, DEU, G, TS, ..., identification of echoes
 measurement procedures:
 - measurement procedures:
 - excitation by impulses (pistols, balloons, ...)
 - excitation by loudspeaker and correlation evaluation (MLS, sinus sweep, ...)
- sensitive with respect to source directivity

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- ▶ general:
 - hints regarding the presence of audience
 - required precision for air temperature measurement: 1°, for humidity measurement 5%
 - sound source: omnidirectional radiation characteristics (specified)
 microphone:
 - microphone.
 - 1/2", pressure microphone
 - (figure of eight microphone for lateral energy fraction)

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instruction for the measurement of reverberation times and further objective criteria:

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measurement geometry of microphone:

- several positions have to be used (representative for the specific conditions)
- minimum distance between two positions: 2 m
- minimum distance to walls: 1 m
- minimum distance to the source: $2\sqrt{\frac{V}{cT}}$ (approx. twice the critical distance)
- height: 1.2 m (height of the ear of a seated person)
- common: 6 microphone positions for smaller rooms (500 seats) and 10 positions for large rooms (2000 seats)

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measurement geometry of the source:

at least 2 positions (representative for activities on stage)
 typical source height: 1.5 m

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dodecahedron loudspeaker for omnidirectional radiation

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"vase"- type loudspeaker for omnidirectional radiation

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example of a measured impulse response



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impulse response measurements with linear microphone array

- vertical display of impulse responses
- wave fronts allow for an estimate of propagation direction



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